in the unit cell, calculated with the values of parameters listed in Table 1, are given in Table 3.

Table 3. Interatomic distances in $\mathrm{HgSO}_{4} . \mathrm{H}_{2} \mathrm{O}$

| Atom | Neighbour | Coordination number | Interatomic distances |
| :---: | :---: | :---: | :---: |
| Hg | O (II) | 2 | $2 \cdot 44 \AA$ |
|  | O (II) | 2 | $2 \cdot 47$ |
|  | $\mathrm{O}(\mathrm{III})$ | 1 | $2 \cdot 10$ |
|  | $\mathrm{O}\left(\mathrm{H}_{2} \mathrm{O}\right)$ | 1 | $2 \cdot 34$ |
| S | O (I) | 1 | 1.24 |
|  | $\mathrm{O}(\mathrm{II})$ | 2 | 1.53 |
|  | O (III) | 1 | 1.53 |
|  | O-O distances in tetrahedron |  |  |
| O(I) | O(II) | - | 1.95 |
|  | O (III) | - | $2 \cdot 24$ |
| O(II) | O (II) | - | 2.53 |
|  | O(III) | - | 2.51 |
|  | $\mathrm{O}-\mathrm{O}$ distances in octahedron |  |  |
| O(II) | O (II) | - | 3.91 |
|  | $\mathrm{O}(\mathrm{II})$ | - | $2 \cdot 89$ |
|  | $\mathrm{O}\left(\mathrm{H}_{2} \mathrm{O}\right)$ | - | $3 \cdot 09$ |
|  | $\mathrm{O}\left(\mathrm{H}_{2} \mathrm{O}\right)$ | - | 3.18 |
|  | O (III) | - | $3 \cdot 85$ |
|  | O (III) | - | $3 \cdot 14$ |

The values for the distances from mercury to the sulphate-oxygen atoms O (II) ( $2 \cdot 44$ and $2.47 \AA$ ) are nearly equal to the sum of ionic radii ( $2 \cdot 45 \AA$ ), if the values of $1 \cdot 10 \AA$, the generally accepted value for the radius of the mercuric ion, and $1.35 \AA$ for oxygen are
used. Distances of nearly the same length ( 2.40 and $2 \cdot 42 \AA$ ) are found to exist also in crystals of mercurous nitrate dihydrate, between mercury and the nitrateoxygen atoms (Grdenić, 1956).

Shorter distances between the mercury atom and the water molecule ( $2 \cdot 34 \AA$ ) are found. It may be compared with the distances between the mercury atom and the water molecule in mercurous nitrate dihydrate ( $2 \cdot 15 \AA$ ). A short coordination $\mathrm{Hg}-\mathrm{O}$ bond has been shown to exist also in mercury diethylene oxide ( $2 \cdot 21 \AA$ ) by Grdenić (1952), and much shorter ( $2 \cdot 03 \AA$ ) in trichlormercury oxonium chloride (Śćavničar \& Grdenić, 1955) and in mercuric oxide (Aurivilius, 1954).

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# The Crystal Structure of $\mathbf{U R e}_{2}$ 

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The intermetallic compound $\mathrm{URe}_{2}$ has two allotropic forms. Up to $180^{\circ} \mathrm{C}$. it has an orthorhombic structure, space group Cmca; above this temperature a C14 structure is formed by a simple dilation of the unit cell. The lattice parameters at $20^{\circ} \mathrm{C}$. are

$$
a=5 \cdot 600, b=9 \cdot 180, c=8 \cdot 460 \AA .
$$

For the hexagonal modification at $213^{\circ} \mathrm{C} . a=5 \cdot 433, c=8.561 \AA$. In both modifications the thermal expansion coefficients are positive along $c$ and negative along $a$.

## Introduction

During an investigation of the uranium-rhenium phase diagram an intermetallic compound $\mathrm{URe}_{2}$ was observed (Brook \& Hatt, 1954). An alloy of composition uranium $66.7 \mathrm{at} . \%$ rhenium appeared single-phased when examined microscopically and by the X-ray Debye-Scherrer powder method, whereas alloys with
higher rhenium content showed free rhenium, and those of lower rhenium content free uranium. Spacing measurements indicated a negligible solubility of either rhenium or uranium in URe 2 .

## Structure determination

Alloys of $\mathrm{URe}_{2}$ were prepared from uranium bar
supplied by the Atomic Energy Research Establishment, Harwell, and rhenium sponge of $99.9 \%$ purity. The sponge rhenium was first arc-melted into buttons on a water-cooled copper hearth under a purified argon atmosphere. These were then arc-melted with uranium to give alloy buttons of approximately 10 g . Heattreatments were carried out by sealing the buttons contained in an alumina crucible inside an evacuated silica glass capsule. No chemical analyses have been made but the melting losses were less than $0.5 \%$.

Debye-Scherrer powder photographs were taken from the crushed alloy, after a heat-treatment of two weeks at $1000^{\circ} \mathrm{C}$. followed by water quenching. The alloy was very brittle so that crushing introduced negligible cold work. A powder pattern taken with Co $K \alpha$ radiation showed 58 resolvable diffractions. These could not be indexed and so rotation singlecrystal photographs were taken using an irregular shaped fragment, selected from the crushed alloy, which showed a pseudo six-fold axis of symmetry. Subsequently this was shown to consist of several grains. Attempts to obtain single grains were unsuccessful. The diffractions on the rotation photograph about one of the axes of two-fold symmetry did not lie on straight layer lines but were scattered about possible layer-line positions, indicating a repeat distance of $5 \cdot 6 \AA$. The rotation photograph about the pseudo six-fold symmetry axis showed straight layer lines giving a repeat distance of $8.2 \AA$. Using these two values all the diffractions on this latter film could be indexed with an orthorhombic cell with dimensions $5 \cdot 6,9 \cdot 2$ and $8 \cdot 2 \AA$. These cell constants enabled all the powder diffractions to be indexed, and the follow-


Fig. 1. Projection of the atoms of a $\mathrm{URe}_{2} \mathrm{Cl} 4$ structure on to a (0001) plane. The letters refer to the interatomic distances given in Fig. 3.
ing systematic absences were observed: 1) $h k l$ when $h+k=2 n+1$, 2) $h 0 l$ when $l=2 n+1$. These conditions are satisfied by space groups $\mathrm{Cmcm}-D_{2 h}^{17}$ and $\mathrm{Cmc} 2_{1}-\mathrm{C}_{2 v}^{12}$. Oscillation photographs taken round the pseudo sixfold axis showed that the fragment consisted of several grains, all with a common $8.2 \AA$ axis but oriented round this axis at approximately $120^{\circ}$ intervals.

Uranium compounds of the form $\mathrm{U} X_{2}$ often crystallize with structures of the so-called Laves phases. These tend to form when the ratio of atomic diameters of the elements is about $1 \cdot 2$. For U and Re this ratio is $1 \cdot 16$ and so its structure was expected to resemble one of the Laves phases. In order to determine the atomic sites of the present orthorhombic structure its relation to the Cl4, hexagonal, Laves structure was considered.

Fig. 1 shows the projection of a Cl4 structure on to a (0001) plane together with the orthohexagonal cell derived from it. If the atomic positions are interpreted in terms of the orthorhombic space group Cmcm then these become, with respect to the equivalent positions ( $0,0,0 ; \frac{1}{2}, \frac{1}{2}, 0$ ),

$$
\begin{aligned}
& 8 \mathrm{U} \text { atoms in sites }(f) \\
& \quad 0, y_{1}, z_{1} ; 0, \bar{y}_{1}, \bar{z}_{1} ; 0, y_{1}, \frac{1}{2}-z_{1} ; 0, \bar{y}, \frac{1}{2}+z_{1} \\
& 8 \operatorname{Re} \text { atoms in sites }(g) \\
& \quad x_{2}, y_{2}, \frac{1}{4} ; \bar{x}_{2}, y_{2}, \frac{1}{4} ; x_{2}, \bar{y}_{2}, \frac{3}{4} ; \bar{x}_{2}, \bar{y}_{2}, \frac{3}{4} \\
& 4 \operatorname{Re} \text { atoms in sites }(c) \\
& \quad 0, y_{3}, \frac{1}{4} ; 0, \bar{y}_{3}, \frac{3}{4} \\
& 4 \operatorname{Re} \text { atoms in sites (a) } \\
& \quad 0,0,0 ; 0,0, \frac{1}{2} .
\end{aligned}
$$

Since the Cl4 structures have axial ratios about 1.63 then the ideal orthohexagonal cell has $a: b: c=$ $1: \sqrt{3}: 1 \cdot 63$. The actual orthorhombic cell determined for $\mathrm{URe}_{2}$ has $a: b: c=1: \gamma 2 \cdot 7: 1 \cdot 52$, and since $a: b$ is no longer exactly $\gamma / 3$ the hexagonal diffractions split into groups with differing multiplicities as shown in Table 1.

Table 1. Comparison of the hexagonal diffractions with those for the orthorhombic cell when b/a deviates from $1 / 3$

| Hexagonal |  | Orthohexagonal (distorted) |  |
| :--- | ---: | :---: | :---: |
| $h k i l$ | $p$ |  | Number of |
| diffractions | $P$ |  |  |
| $0,0,0, l$ | 2 | Single | 2 |
| $h, h, \overline{2} h, 0$ | 6 | Double | 2,4 |
| $h, h, \overline{2} h, l$ | 12 | Double | 4,8 |
| $0, k, \bar{k}, 0$ | 6 | Double | 2,4 |
| $0, k, \bar{k}, l$ | 12 | Double | 4,8 |
| $h, k, i, 0$ | 12 | Triple | $4,4,4$ |
| $h, k, i, l$ | 24 | Triple | $8,8,8$ |

The intensities were calculated from the formula $I \propto p\left|F_{h k l}\right|^{2}\left(1+\cos ^{2} 2 \theta\right) / \sin ^{2} \theta \cos \theta$, in the usual notation. Satisfactory agreement with observed values were obtained with $y_{1}=\frac{1}{3}, z_{1}=\frac{1}{16}, x_{2}=\frac{1}{4}, y_{2}=\frac{1}{12}, y_{3}=-\frac{1}{6}$. No detailed refinement of these parameters has been undertaken but they are probably correct to $\pm 0.01$. Observed and calculated values of intensities and $d$ spacings are given in Table 2.

Table 2. Observed and calculated values of d spacings and intensities for the orthorhombic modification of $\mathrm{URe}_{2}$ at $20^{\circ} \mathrm{C}$.

Wavelength of X-radiation used Co $K \alpha_{1}, \lambda=1.78890 \AA$

$v w=$ very weak ; $w=$ weak; $m=$ medium; $s=$ strong; $b=$ broad.

## High temperature modification

The room-temperature structural work indicated the possibility that the orthorhombic form of $\mathrm{URe}_{2}$ had transformed from a hexagonal C14 structure while
cooling. To check this prediction several photographs were taken with a Unicam high-temperature camera. The furnace thermo-couples were not calibrated throughout the range of temperature investigated, but exposures made with a super-purity aluminium

Table 3. Observed and calculated values of $d$ spacings and intensities for the C14 hexagonal modification of $\mathrm{URe}_{2}$ at $655^{\circ} \mathrm{C}$.
Wavelength of X-radiation used $\mathrm{Cu} K \alpha_{1}, \lambda=1 \cdot 54050 \AA$.
 perature range 100 to $200{ }^{\circ} \mathrm{C}$. Temperatures in this range have been corrected for this error and are probably correct to $\pm 10^{\circ}$ whereas above $200^{\circ} \mathrm{C}$. the same error has been assumed but the accuracy will not be as great.
The $\mathrm{URe}_{2}$ powder was sealed in a silica quill after evacuating to a pressure of the order $10^{-2} \mathrm{~mm} . \mathrm{Hg}$. X-ray photographs were taken at $20,151,210,340,555$ and $655^{\circ} \mathrm{C}$. Before each exposure the powder was maintained at the temperature for 2 hr . Traces of $\mathrm{UO}_{2}$ appeared after 5 hr . at $655^{\circ} \mathrm{C}$.

Up to $151{ }^{\circ} \mathrm{C}$. the diffractions corresponded to the orthorhombic version of $\mathrm{URe}_{2}$ already described. At $210^{\circ} \mathrm{C}$. and above a new set of diffractions appeared which could all be indexed on a hexagonal cell. The systematic absences and line intensities are consistent with the space group and atomic sites of the C14 hexagonal Laves phase, as given in International Tables for X-ray Crystallography (1952).

Space group $P 6_{3} / m m c-D_{6 k}^{4}$.
4 U atoms in sites $(f)$

$$
\begin{aligned}
& \frac{1}{3}, \frac{2}{3}, z ; \frac{2}{3}, \frac{1}{3}, \bar{z} ; \frac{2}{3}, \frac{1}{3}, \frac{1}{2}+z ; \frac{1}{3}, \frac{2}{3}, \frac{1}{2}-z \\
& \text { with } z=\frac{1}{16} \pm 0.01
\end{aligned}
$$

6 Re atoms in sites ( $h$ )

$$
\begin{aligned}
& x, 2 x, \frac{1}{4} ; 2 \bar{x}, \bar{x}, \frac{1}{4} ; x, \bar{x}, \frac{1}{4} ; \bar{x}, 2 \bar{x}, \frac{3}{4} \\
& 2 x, x, \frac{3}{4} ; \bar{x}, x, \frac{3}{4} \\
& \quad \text { with } x=-\frac{1}{6} \pm 0.01
\end{aligned}
$$

The diffraction data are listed in Table 3 and show good agreement between observed and calculated values of intensities and $d$ spacings. Lattice parameters of both modifications and their variation with temperature are given in Table 4.

Table 4. Cell dimensions at various temperatures

|  | Temp. ( ${ }^{\circ} \mathrm{C}$.) | Lattice parameter ( $\AA$ ) |  |  | Cell Vol. |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $a$ | $b$ | c |  |
| Orthorhombic | 20 | 5.600 | 9-178 | $8 \cdot 463$ | 435 |
|  | 151 | 5.555 | 9•180 | $8 \cdot 512$ | 434 |
| Hexagonal |  | $a$ | $a / 3$ | $c$ |  |
|  | 213 | 5.433 | $9 \cdot 410$ | $8 \cdot 561$ | 438 |
|  | 340 | $5 \cdot 420$ | $9 \cdot 389$ | 8.590 | 437 |
|  | 655 | $5 \cdot 405$ | 9.363 | 8.682 | 441 |

The above cell parameters are correct to $\pm 0.003 \AA$.

## Discussion

The hexagonal high-temperature modification of $\mathrm{URe}_{2}$ transforms on cooling to an orthorhombic structure, at a temperature which lies between 210 and $151^{\circ} \mathrm{C}$. with a $2.0 \%$ expansion in the $a$ direction and contractions of 2.4 and $0.35 \%$ in the $b$ and $c$ directions respectively. The original hexagonal grain transforms
into several orthorhombic domains each having one of the three orthohexagonal orientations, thus producing a macroscopic pseudo hexagonal grain. The transformation may be termed martensitic since it consists of a simple dilation of the unit cells without breaking the atomic bonds.


Fig. 2. Variation of cell dimensions and volume with temperature.

The anisotropic thermal expansion of both modifications are shown in Fig. 2 and the variation in near neighbour atomic distances in Fig. 3. The latter shows that the extrapolated curves intersect at approximately $1000^{\circ} \mathrm{C}$. when all $\mathrm{U}-\mathrm{U}$ distances are equal to $3 \cdot 29$, Re-Re $2 \cdot 69$ and U-Re $3 \cdot 15 \AA$. At this temperature the Cl 4 structure would have the ideal axial ratio of $1 \cdot 63$. As the temperature falls, the distance between planes of atoms perpendicular to the $c$ direction decreases, whereas distances in these planes increase. During the transformation those interatomic distances which contracted while cooling undergo small changes in length but those which expanded undergo relatively large ones. This suggests that the origin of the transformation is that on cooling the strain in the $\{000 l\}$


Fig. 3. Variation of interatomic distances with temperature. The interatomic distance $c$ is that between uranium atoms vertically above one another; the others are as shown in Fig. 1.
planes produced by the thermal expansion reaches a maximum value at $\sim 180^{\circ} \mathrm{C}$. and is relieved by disrupting the $\{000 l\}$ hexagonal atomic arrangement. Conversely, on heating, the atoms have sufficient energy at $180^{\circ} \mathrm{C}$. to take on the more symmetrical arrangement.

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